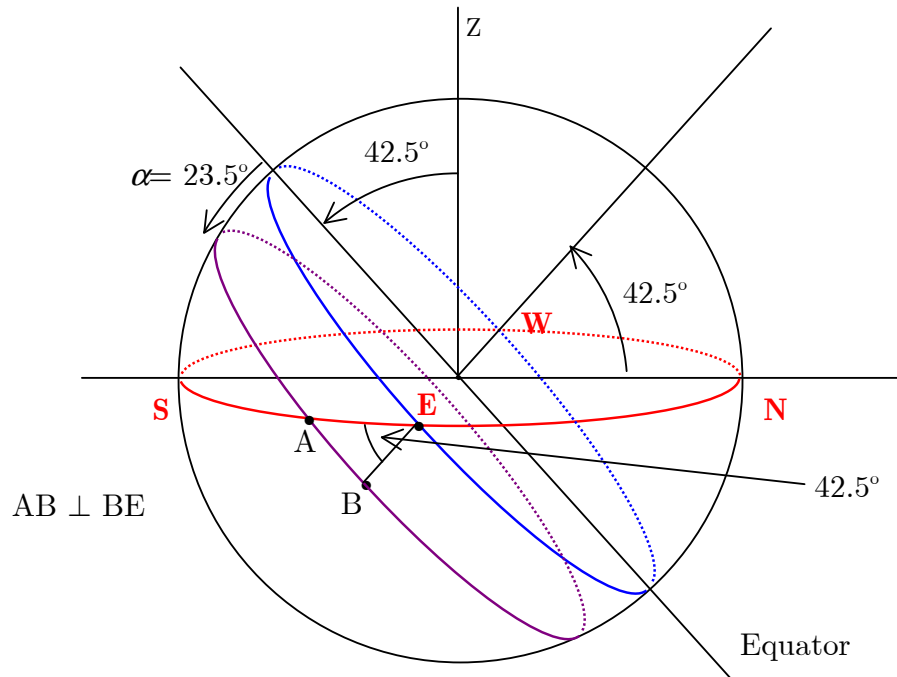


**SOLUTION FOR QUESTION 1. (30 points for 15 short questions)**

1.1



E = East point

$$BE = 23.5^\circ = |\text{declination of the Sun}|$$

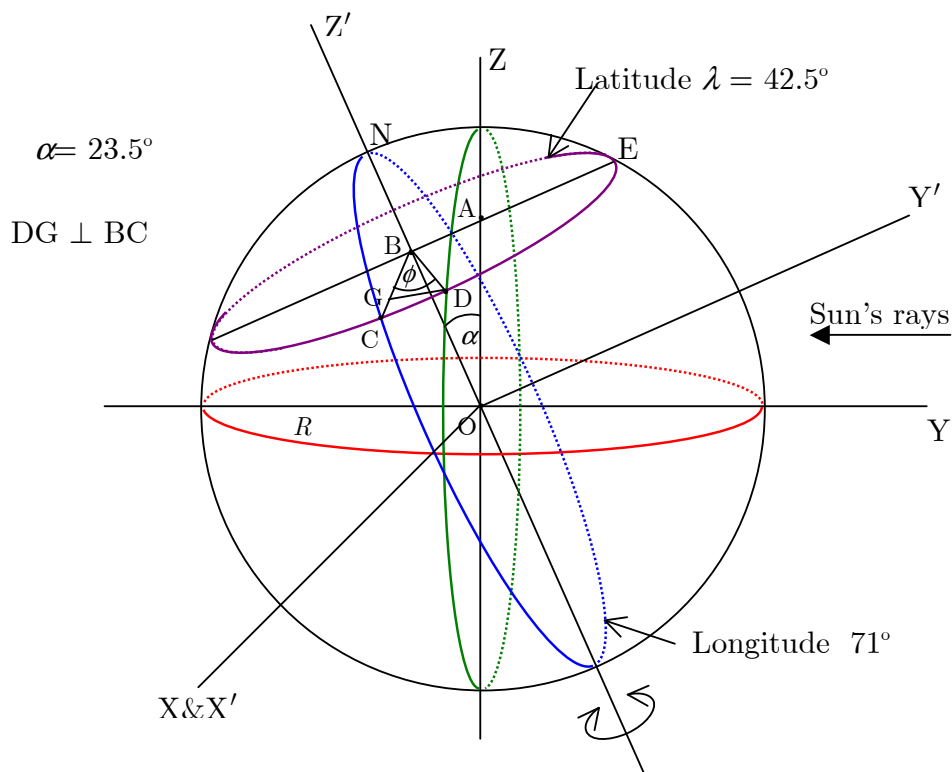
$$\frac{AB}{\sin 42.5^\circ} = \frac{BE}{\sin(90^\circ - 42.5^\circ)}$$

$$AB = BE \frac{\sin 42.5^\circ}{\cos 42.5^\circ} = 23.5^\circ \tan 42.5^\circ$$

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local time =  $(23.5 \tan 42.5^\circ / 15)$  hrs after 6:00 = 7:26 am. The official time at  $75^\circ$  W should be 16 min. less.

**Ans. 7:10 am.**

Alternative solution (for 1.1)

The Earth's position relative to the Sun is shown in the figure.

Note that  $OB = R \sin \lambda$

$$AB = OB \tan \alpha$$

$$BC = BD = BE = R \cos \lambda$$

$$\sin \phi = \frac{DG}{BD} = \frac{BA}{BD} = \frac{BO \tan \alpha}{R \cos \lambda} = \tan \alpha \tan \lambda$$

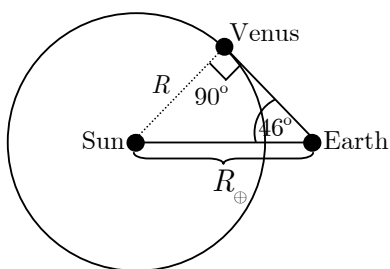
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$$= \tan(23.5^\circ) \tan(42.5^\circ) = 0.39843 = \sin(23.48^\circ)$$

Hence, the Sun will rise at  $(71^\circ + 23.48^\circ) \times 4 \text{ min.} - 5 \text{ hours.} = 77.92 \text{ min.}$

after 6 a.m.. This is at 7:18 a.m.

1.2



The angular separation is maximum when Sun, Venus and Earth form a right-angled triangle as shown.

$$\begin{aligned}
 \text{Here } R &= R_{\oplus} \sin 46^\circ \\
 &= (1 \text{ A.U.}) \sin 46^\circ \\
 &= 0.72 \text{ A.U.}
 \end{aligned}$$

1.3 If the same face of the Earth were to face the Sun all the time then the Earth would make one complete turn relative to fixed stars in one solar year (365.25 solar days).

This implies that in 365.25 solar days our actual Earth makes (365.25+1) complete turns relative to fixed stars.

Hence 365.25 solar days are the same time interval as 366.25 sidereal days;

$$\begin{aligned}
 \text{and } 183 \text{ solar days} &\equiv \frac{183 \times 366.25}{365.25} \text{ sidereal days} \\
 &= 183.50 \text{ sidereal days}
 \end{aligned}$$

**OR**

$$1 \text{ solar day} = 24 \times 60 \times 60 \text{ s} = 86400 \text{ s}$$

$$1 \text{ sidereal day} = 23 \times 3600 + 56 \times 60 + 4.1 = 86164.1 \text{ s}$$

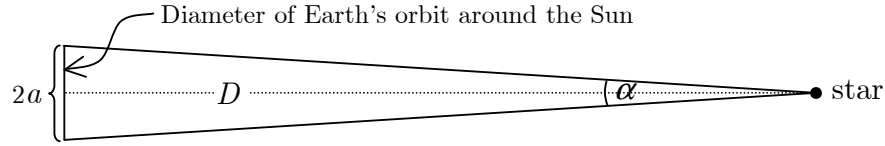
$$183 \text{ solar days} = 183.50 \text{ sidereal days}$$

1.4 During a full Moon we see the whole face of the Moon.

$$\begin{aligned}
 \text{Hence } \frac{\text{Moon's diameter}}{\text{distance to Moon}} &= \text{angle in radians} \\
 &= \frac{0.46 \times \pi}{180}
 \end{aligned}$$

$$\begin{aligned}
 \text{Distance to the Moon} &= (\text{Moon's diameter}) \times \frac{180}{0.46 \times \pi} \\
 &= (2 \times 1.7374 \times 10^6 \text{ m}) \times \frac{180}{0.46 \times \pi} \\
 &= 4.328 \times 10^8 \text{ m} = 4.3 \times 10^5 \text{ km}
 \end{aligned}$$

1.5



$$a = 1 \text{ A.U.} = 1.496 \times 10^{11} \text{ m}$$

$$1 \text{ pc} = 3.0856 \times 10^{16} \text{ m}$$

$$D = 100 \times 3.0856 \times 10^{16} \text{ m}$$

$$\begin{aligned} \alpha &= \frac{2a}{D} = \frac{2 \times 1.496 \times 10^{11}}{100 \times 3.0856 \times 10^{16}} \text{ radian} \\ &= 0.96966 \dots \times 10^{-7} \text{ radian} \\ &= 5.555779 \dots \times 10^{-6} \text{ degree} \\ &= 0.02 \text{ arc second} \end{aligned}$$

1.6 According to Kepler's third law we have

$$\begin{aligned} (\text{period})^2 &= (\text{constant})(\text{semi-major axis})^3 \\ T^2 &= (\text{constant})a^3 \end{aligned}$$

This constant is  $1 \frac{(\text{year})^2}{(\text{A.U.})^3}$  when  $T$  is measured in years and  $a$  in A.U.'s.

For this comet we have

$$\begin{aligned} a &= \frac{31.5 + 0.5}{2} = 16.0 \text{ A.U.} \\ T^2 &= (16)^3 = 16 \times 4 \times 4 \times 16 = (64)^2 \\ T &= 64 \text{ years} \end{aligned}$$

1.7 According to Kepler's second law we have;

The area is swept out at constant rate throughout the orbital motion.

Now, for this comet the area swept out in one orbital period is  $\pi ab$  where  $a$  is the semi-major axis and  $b$  the semi-minor axis of the orbit.

$$a = 16.0 \text{ A.U.}$$

$$b^2 = a^2 - (\text{distance from a focus to centre of ellipse})^2$$

$$= (16.0)^2 - (16.0 - 0.5)^2$$

$$b = 3.968 \text{ A.U.}$$

$$\pi ab = 199.5 \text{ (A.U.)}^2$$

Hence, the area swept out per year is

$$\frac{199.5 \text{ (A.U.)}^2}{64 \text{ years}} = 3.1 \text{ (A.U.)}^2 / \text{year}$$

1.8 From the Wien's displacement law

$$\lambda_{\max} T = 2.8977 \times 10^{-3} \text{ m.K}$$

$$\lambda_{\max} = \frac{2.8977 \times 10^{-3}}{4000} \text{ m} = 7.244 \times 10^{-7} \text{ m}$$

$$= 724 \text{ nanometers}$$

1.9 From the Stefan-Boltzmann law, the total power emitted is

$$L = (4\pi R^2)(\sigma T^4) = (4\pi R_{\odot}^2)(\sigma T_{\odot}^4) \left( \frac{R}{R_{\odot}} \right) \left( \frac{T}{T_{\odot}} \right)^4 = L_{\odot} (2.5)^2 \left( \frac{7500}{5800} \right)^4$$

$$\frac{L}{L_{\odot}} = (6.25)(2.796) = 17.47 = 17.5$$

$$L = 17.5 L_{\odot}$$

1.10 flux density =  $\frac{\text{luminosity}}{4\pi \times (\text{distance})^2}$

$$\text{distance} = \sqrt{\frac{0.4 \times 3.826 \times 10^{26}}{4\pi \times 6.23 \times 10^{-14}}} = 1.398 \times 10^{19} \text{ m}$$

$$= \frac{1.398 \times 10^{19}}{3.0856 \times 10^{16}} \text{ pc} = 453 \text{ pc}$$

1.11  $L_{\text{supernova}} = 10^{10} L_{\odot}$

Let  $D$  be the distance to that supernova.

Then the flux intensity on Earth would be  $I = \frac{10^{10} L_{\odot}}{4\pi D^2}$

This must be the same as  $\frac{L_{\odot}}{4\pi (D_{\oplus-\odot})^2}$

$$\frac{10^{10} L_{\odot}}{4\pi D^2} = \frac{L_{\odot}}{4\pi (1 \text{ A.U.})^2}$$

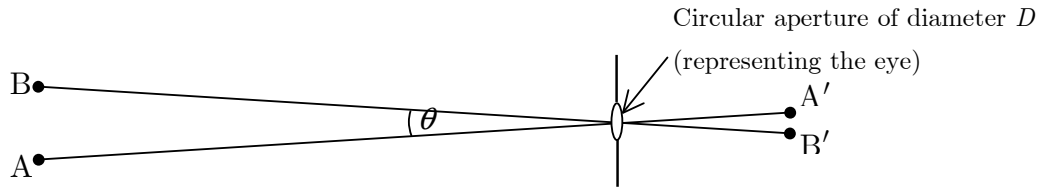
$$\begin{aligned} \therefore D &= 10^{\frac{10}{2}} \text{ A.U.} = 10^5 \text{ A.U.} \\ &= \frac{1.496 \times 10^{16}}{3.0856 \times 10^{16}} \text{ pc} = 0.485 \text{ pc} \\ &= 0.485 \times 3.2615 \text{ ly} = 1.58 \text{ ly} \end{aligned}$$

- 1.12 The difference in frequencies is due to the relativistic Doppler shift. Since the observed frequency of emission from the gas cloud is higher than the laboratory frequency  $\nu_0$ , the gas cloud must be approaching the observer.

$$\begin{aligned} \text{Hence, } \nu &= \nu_0 \sqrt{\frac{c+v}{c-v}} \\ \frac{v}{c} &= \frac{\left(\frac{\nu}{\nu_0}\right)^2 - 1}{\left(\frac{\nu}{\nu_0}\right)^2 + 1} \\ &= \frac{0.001752379}{2.001752379} = 0.000875422 \\ v &= (0.000875422)(2.99792458 \times 10^8 \text{ m.s}^{-1}) \\ &= 0.00262445 \times 10^8 \text{ m.s}^{-1} \\ &= 262.445 \text{ km.s}^{-1} \end{aligned}$$



1.13



A and B are two distant objects.

A' and B' are the central maxima of their diffracted images.

The angular position  $\phi_1$  of the first minimum relative to central maximum

of each diffraction pattern is given by  $\phi_1 = 1.22 \frac{\lambda}{D}$ .

According to Lord Rayleigh the minimum angle of resolution is  $\phi_1$ .

Hence the images of A and B will be resolved if  $\theta > \phi_1$ ,  $D > 1.22 \frac{\lambda}{\theta}$ .

We may take distance AB to be the diameter of the crater. Hence the

diameter  $D$  of the eye's aperture must be, at least,  $D_{\min} = 1.22 \frac{\lambda}{\theta}$ .

$$\begin{aligned} \theta &= \frac{80 \text{ km}}{\text{distance from Earth to Moon}} \\ &= \frac{80 \times 10^3 \text{ m}}{3.844 \times 10^8 \text{ m}} = 2.081 \times 10^{-4} \text{ radian} \end{aligned}$$

$$\lambda = 500 \times 10^{-9} \text{ m}$$

$$\begin{aligned} \therefore D_{\min} &= \frac{1.22 \times 500 \times 10^{-9}}{2.081 \times 10^{-4}} \text{ m} = 2.93 \times 10^{-3} \text{ m} \\ &= 2.9 \text{ mm} \end{aligned}$$

Hence, it is possible to resolve the 80 km-diameter crater with naked eye.

1.14 The spherical boundary at which the escape velocity becomes equal to the

speed of light is of radius  $R = \frac{2GM_{\odot}}{c^2}$

$$R = \frac{2 \times 6.672 \times 10^{-11} \times 1.989 \times 10^{30}}{(2.99792458 \times 10^8)^2} \text{ m}$$

$$\begin{aligned}
&= \frac{2 \times 6.672 \times 1.989}{2.998 \times 2.998} \times 10^3 \text{ m} \\
&= 2.95 \text{ km}
\end{aligned}$$

1.15 The flux ratio versus magnitude difference implies

$$\begin{aligned}
m_1 - m_2 &= -2.5 \log \left( \frac{f_1}{f_2} \right) \\
\frac{f_1}{f_2} &= 10^{\frac{(m_2 - m_1)}{2.5}}
\end{aligned}$$

So for a magnitude difference of  $(-1.5) - 6 = -7.5$  we find a flux ratio of

$$\frac{f_{\min}}{f_{\max}} = 10^{\frac{-7.5}{2.5}} = 10^{-3}$$

And thus the visible stars range only over a factor of 1000 in brightness.